



LUND UNIVERSITY



For tutorials  
at Summer Schools

# PYTHIA 8 Worksheet

Torbjörn Sjöstrand

Department of Theoretical Physics, Lund University

Peter Skands

Theoretical Physics Department, Fermilab

## 1 Introduction

The objective of this exercise is to teach you the basics of how to use the PYTHIA 8.1 event generator to study various physics aspects. As you become more familiar you will better understand the tools at your disposal, and can develop your own style to use them. Within this first exercise it is not possible to describe the physics models used in the program; for this we refer to the PYTHIA 8.1 brief introduction [1], to the full PYTHIA 6.4 physics description [2], and to all the further references found in them.

PYTHIA 8 is, by today's standards, a small package. It is completely selfcontained, and is therefore easy to install for standalone usage, e.g. if you want to have it on your own laptop, or if you want to explore physics or debug code without any danger of destructive interference between different libraries. Section 2 describes the procedure, which is what we will need for this introductory session.

When you use PYTHIA you are expected to write the main program yourself, for maximal flexibility and power. Several examples of such main programs are included with the code, to illustrate common tasks and help getting started. Section 3 gives you a simple step-by-step recipe how to write a minimal main program, that could then gradually be expanded in different directions, e.g. as in Section 4.

In Section 5 you will see how the parameters of a run can be read in from a file, so that the main program can be kept fixed. Many of the provided main programs therefore allow you to create executables that can be used for different physics studies without recompilation, but potentially at the cost of some flexibility.

While PYTHIA can be run standalone, it can also be interfaced with a set of other libraries. One example is HEPMC, which is the standard format used by experimentalists to store generated events. Since the HEPMC library location is installation-dependent it is not possible to give a fool-proof linking procedure, but some hints are provided for the interested in Section 6.

Finally Section 7 gives some suggestions for the study of other possible physics topics, and Appendix A a brief summary of the event-record structure.

## 2 Installation

Denoting a generic PYTHIA 8 version `pythia81xx`, here is how to install PYTHIA 8 on a Linux/Unix/MacOSX system as a standalone package.

1. In a browser, go to  
`http://www.thep.lu.se/~torbjorn/Pythia.html`
2. Download the (current) program package  
`pythia81xx.tgz`  
to a directory of your choice (e.g., by right-clicking on the link — the left-click default location is installation-dependent, and could be your home directory or the Desktop one.)
3. In a terminal window, `cd` to where `pythia81xx.tgz` was downloaded, and type  
`tar xvfz pythia81xx.tgz`  
This will create a new (sub)directory `pythia81xx` where all the PYTHIA source files are now ready and unpacked.
4. Move to this directory (`cd pythia81xx`) and do a `make`. This will take 1–3 minutes (computer-dependent). The PYTHIA 8 libraries are now compiled and ready for physics.
5. For test runs, `cd` to the `examples/` subdirectory. An `ls` reveals a list of programs, `mainNN`, with `NN` from 01 through 27. These example main programs each illustrate an aspect of PYTHIA 8. For a list of what they do, see the `README` file in the same directory.

Initially only use one or two of them to check that the installation works. Once you have worked your way through the introductory exercises in the next sections you can return and study the programs and their output in more detail.

To execute one of the test programs, do

```
make mainNN
./mainNN.exe
```

The output is now just written to the terminal, `stdout`. To save the output to a file instead, do `./mainNN.exe > mainNN.out`, after which you can study the test output at leisure by opening `mainNN.out`. See Appendix A for an explanation of the event record that is listed in several of the runs.

6. If you open the file  
`pythia81xx/html/doc/Welcome.html`  
you will gain access to the online manual, where all available methods and parameters are described. Use the left-column index to navigate among the topics, which then are displayed in the larger right-hand field.

## 3 A “Hello World” program

We will now generate a single  $gg \rightarrow t\bar{t}$  event at the LHC, using PYTHIA standalone.

Open a new file `mymain.cc` in the `examples` subdirectory with a text editor, e.g. Emacs. Then type the following lines (here with explanatory comments added):

```
// Headers and Namespaces.
#include "Pythia.h" // Include Pythia headers.
using namespace Pythia8; // Let Pythia8:: be implicit.

int main() { // Begin main program.

    // Set up generation.
    Pythia pythia; // Declare Pythia object
    pythia.readString("Top:gg2ttbar = on"); // Switch on process.
    pythia.init( 2212, 2212, 14000.); // initialize pp (PDG 2212) at LHC.

    // Generate event(s).
    pythia.next(); // Generate an(other) event. Fill event record.
    pythia.event.list(); // Print contents of event record.

    return 0; } // End main program with error-free return.
```

Next you need to edit the `Makefile` (the one in the `examples` subdirectory) to make it know what to do with `mymain.cc`. The lines

```
# Create an executable for one of the normal test programs
main00 main01 main02 main03 ... main09 main10 main10 \
```

and the two next enumerate the main programs that do not need any external libraries. Edit the last of these lines to include also `mymain`:

```
main21 main22 main23 main24 main25 main26 main27 mymain: \
```

Now it should work as before with the other examples:

```
make mymain
./mymain.exe > mymain.out
```

whereafter you can study `mymain.out`, especially the example of an event record, see Appendix A. For instance, see if you can find several copies of the top quarks, and check the status codes to figure out why each new copy has been added. Also note how the mother/daughter indices tie the various copies together.

## 4 A first realistic analysis

We will now gradually expand the skeleton main program from above, towards what would be needed for a more realistic analysis setup.

- Often, we wish to mix several processes together. To add in the process  $q\bar{q} \rightarrow t\bar{t}$  to the above example, just include a second `pythia.readString` call with argument `"Top:qqbar2ttbar = on"`.
- To give some variety of event types, introduce a loop over `pythia.next()`, so that it is called a few times (say 5 or 10 to begin with). Each call resets the event record and fills it with a new event.

- To obtain statistics on the number of events generated of the different kinds, and the estimated cross sections, add a  

```
pythia.statistics()
```

just before the end of the program.

- During the run you may receive problem messages. These come in three kinds:
  - a *warning* is a minor problem that is easily fixed, at least approximately;
  - an *error* is a bigger problem, that can normally still be fixed by backing up and trying again;
  - an *abort* is such a major problem that the current event could not be completed; in such a rare case `pythia.next()` is `false` and the event should be skipped.

A message is printed the first time for each new problem kind, while the above-mentioned `pythia.statistics()` tells how many times each problem was encountered.

- Looking at the `pythia.event.list()` listing for a few events at the beginning of each run is useful to make sure you are generating the right kind of events, at the right energies, etc. But for real analyses, you need automated access to the event record. The PYTHIA event record provides many utilities to make this as simple and efficient as possible. To loop over all particles in the event record, use

```
for (int i = 0; i < pythia.event.size(); ++i) {
```

with a matching closing curly bracket further down.

- Inside the loop you can access the properties of particle `pythia.event[i]`. For instance, the method `id()` returns the PDG identity code (see Appendix A), so if you insert

```
cout << "i = " << i << " id = " << pythia.event[i].id() << endl;
```

you should get a table of PDG identity codes that matches the one in the event listing.

- The event listing contains all partons and particles, traced through a number of intermediate steps. For instance, several copies of the top may well appear in the listing, as the addition of showers shifts the top around. It is the last copy of top that gives the “final” answer. You can thus obtain the location of this top e.g. by a line

```
int iTop = 0;
```

before the loop over the particles in the event, and

```
if (pythia.event[i].id() == 6) iTop = i;
```

inside of it.

- In addition to the particle properties in the event listing there are also methods that return many derived quantities for a particle, such as transverse momentum, `pythia.event[i].pT()`, and pseudorapidity, `pythia.event[i].eta()`. Use these methods to print out the values for the final top found above.

- We now want to generate more events, say 1000, to view the shape of these distributions. You then need to remove the listing and printing for each event. In its place book two histograms (a very simple histogramming class comes along with the program, for rapid check/debug purposes)

```
Hist pT("top transverse momentum", 100, 0., 200.);
```

```
Hist eta("top pseudorapidity", 100, -5., 5.);
```

before the event loop, store the values

```

    pT.fill( pythia.event[iTop].pT() );
    eta.fill( pythia.event[iTop].eta() );
for each event, and write out the histograms
    cout << pT << eta;
after the event loop.

```

- As a final standalone exercise, consider plotting the charged multiplicity of events. You then need to have a counter set to zero for each new event. Inside the particle loop this counter should be updated whenever the particle `isCharged()` and `isFinal()`. For the histogram, note that it can be treacherous to have bin limits at integers, where roundoff errors decide whichever way they go. In this particular case only even numbers are possible, so 100 bins from  $-1$  to 399 would still be acceptable.

## 5 Input files

With the `mymain.cc` structure developed above it is necessary to recompile the main program for each minor change, e.g. if you want to rerun with more statistics. This is not time-consuming for a simple standalone run, but may become so for more realistic applications. Therefore parameters can be put in a special input “cards” file that is read by the main program.

The name of this input file can be hardcoded in the main program. However, for more flexibility, it can also be provided as a command-line argument. Then replace the `int main() {` line by

```

    int main(int argc, char* argv[]) {
and all the readString(...) commands by a single
    pythia.readFile(argv[1]);

```

The executable `mymain.exe` is then run with a command line like

```
./mymain.exe infile > mymain.out
```

Here `infile` is your chosen name for the input file we will describe in the following.

The `infile` can contain one command per line, of the type

```
variable = value
```

All PYTHIA variable names begin with a letter or a digit, so if you begin a line with a non-alphanumeric character (such as `!`, `#` or `$`) it will be interpreted as a comment line. You can also add comments on a valid command line, after the value. For readability we recommend it to be preceded by a non-alphanumeric character, as for comment lines. Blanks can be freely inserted, except inside the variable name and the value. Further, the equal sign is optional, but improves readability. The variable names are case-insensitive; the mixing of cases has been chosen purely to improve readability.

As before, all valid variables are listed in the online manual, see section 2 point 6. Cut-and-paste of variable names can be used to avoid spelling mistakes.

The same pages also exist in an interactive variant, where you semi-automatically can construct a file with all the command lines you wish to have. This requires that somebody installs the `pythia81xx/phpdoc` directory in a webserver (i.e. like your homepage is). If you lack a local installation you can use the one at

```
http://home.thep.lu.se/~torbjorn/php8108/Welcome.php
```

This is not a commercial-quality product, however, but requires some user discipline. As explained on the “Save Settings” page, you must initially pick a temporary file, then remember to save changes to this file for each new webpage that you change, and finally determine the ultimate location of the file.

To allow beam parameters to be set in the input file, they should be removed from the arguments of `pythia.init(...)`, since else the arguments take precedence. That is, you should edit the call to read `pythia.init()`.

As a replacement for these, and for the previous `pythia.readString(...)` commands, your `infile` should thus contain the lines (with explanatory comments added)

```
Beams:idA = 2212      ! first incoming beam is a 2212, i.e. a proton.
Beams:idB = 2212      ! second beam is also a proton.
Beams:eCM = 14000.    ! the cm energy of collisions.
Top:gg2ttbar = on     ! switch on the process g g -> t tbar.
Top:qqbar2ttbar = on ! switch on the process q qbar -> t tbar.
```

In addition to all the internal `Pythia` commands there exist a few defined in the database but not actually used. These are intended to be useful in the main program, and thus begin with `Main:.` The most basic of those is `Main:numberOfEvents`, which you can use to specify how many events you want to generate. To make this have any effect, you need to read it in the main program, after the `pythia.readFile(...)` command, by a line like

```
int nEvent = pythia.mode("Main:numberOfEvents");
and set up the event loop like
for (int iEvent = 0; iEvent < nEvent; ++iEvent) {
```

The above commands together would provide a minimal `infile`. As you go along you can gradually increase the number of events and also play with further optional aspects, such as:

- `6:m0 = 175.`  
change the top mass, which by default is 171 GeV.
- `PartonLevel:FSR = off`  
switch off final-state radiation.
- `PartonLevel:ISR = off`  
switch off initial-state radiation.
- `PartonLevel:MI = off`  
switch off multiple interactions.
- `Random:setSeed = on`  
`Random:seed = 123456789`  
all runs by default use the same random-number sequence, for reproducibility, but you can pick any number between 1 and 900,000,000 to obtain a unique sequence.

For instance, check the importance of FSR, ISR and MI on the charged multiplicity of events by switching off one component at a time.

The usage of further `Main:` variables is illustrated e.g. in `main03.cc`, and of the possibility to use command-line input files in `main17.cc` and `main32.cc`.

## 6 Interface to HepMC

The standard HEPMC event-record format will be frequently used in subsequent training sessions, with a ready-made installation. However, for the ambitious, here is described how to set up the PYTHIA interface, assuming you already know where HEPMC is installed. Note: the interface to HEPMC version 1 is no longer supported; you must use version 2.

To begin with, you need to go back to the installation procedure of section 2 and insert/re-do some steps.

1. Move back to the main `pythia81xx` directory (`cd ..` if you are in `examples`).
2. Remove the currently compiled version with  
`make clean`
3. Configure the program in preparation for the compilation:  
`./configure --with-hepmc=path`  
where the directory-tree `path` would depend on your local installation.
4. Should `configure` not recognize the version number you can supply that with an optional argument, like  
`./configure --with-hepmc=path --with-hepmcversion=2.03.06`
5. Recompile the program, now including the HEPMC interface, with `make` as before, and move back to the `examples` subdirectory.
6. Do either of  
`source config.csh`  
`source config.sh`  
the former when you use the `csh` or `tcsh` shells, otherwise the latter. (Use `echo $SHELL` if uncertain.)
7. You can now also use the `main31.cc` and `main32.cc` examples to produce HEPMC event files. The latter may be most useful; it presents a slight generalization of the command-line-driven main program you constructed in Section 5. After you built the executable you can run it with  
`./main32.exe infile hepmcfile > main32.out`  
where `hepmcfile` is your chosen name for the output file with HEPMC events.

Note that the above procedure is based on the assumption that you will be running your main programs from the `examples` subdirectory. If not you will have to create your own scripts and/or makefiles to handle the linking. If you have no experience with such tasks then better use existing instructions for your local installation. If you do have such experience then here a short summary what you need to know to get going.

Before you run a PYTHIA program the `PYTHIA8DATA` environment variable needs to be set to point to the `xmldoc` subdirectory where all settings and particle data are stored. If you use the `csh` or `tcsh` shells this means a line like

```
setenv PYTHIA8DATA /path/pythia81xx/xmldoc
```

and else

```
export PYTHIA8DATA=/path/pythia81xx/xmldoc
```

where the correct `path` has to be found by you. Similarly, to use HEPMC, you also have to set or append its location to the `LD_LIBRARY_PATH` (the `DYLD_LIBRARY_PATH` on Mac OSX); the `config.csh` and `config.sh` files generated above well illustrate the code needed to achieve this. Finally, the necessary linking stage can be understood from the

relevant parts of the `examples/Makefile`.

## 7 Further studies

If you have time left, you should take the opportunity to try a few other processes or options. Below are given some examples, but feel free to pick something else that you would be more interested in.

- One popular misconception is that the energy and momentum of a B meson has to be smaller than that of its mother b quark, and similarly for charm. The fallacy is twofold. Firstly, if the b quark is surrounded by nearby colour-connected gluons, the B meson may also pick up some of the momentum of these gluons. Secondly, the concept of smaller momentum is not Lorentz-frame-independent: if the other end of the b colour force field is a parton with a higher momentum (such as a beam remnant) the “drag” of the hadronization process may imply an acceleration in the lab frame (but a deceleration in the beam rest frame).

To study this, simulate b production, e.g. the process `HardQCD:gg2bbbar`. Identify B/B\* mesons that come directly from the hadronization, for simplicity those with status code `-83` or `-84`. In the former case the mother b quark is in the `mother1()` position, in the latter in `mother2()`. (Study a few event listings to see how it works.) Plot the ratio of B to b energy to see what it looks like.

- One of the characteristics of multiple-interactions (MI) models is that they lead to strong long-range correlations, as observed in data. That is, if many hadrons are produced in one rapidity range of an event, then most likely this is an event where many MI’s occurred (and the impact parameter between the two colliding protons was small), and then one may expect a larger activity also at other rapidities.

To study this, select two symmetrically located, one unit wide bins in rapidity (or pseudorapidity), with a variable central separation  $\Delta y$ :  $[\Delta y/2, \Delta y/2 + 1]$  and  $[-\Delta y/2 - 1, -\Delta y/2]$ . For each event you may find  $n_F$  and  $n_B$ , the charged multiplicity in the “forward” and “backward” rapidity bins. Suitable averages over a sample of events then gives the forward–backward correlation coefficient

$$\rho_{FB}(\Delta y) = \frac{\langle n_F n_B \rangle - \langle n_F \rangle \langle n_B \rangle}{\sqrt{(\langle n_F^2 \rangle - \langle n_F \rangle^2)(\langle n_B^2 \rangle - \langle n_B \rangle^2)}} = \frac{\langle n_F n_B \rangle - \langle n_F \rangle^2}{\langle n_F^2 \rangle - \langle n_F \rangle^2},$$

where the last equality holds for symmetric distributions such as in pp and  $\bar{p}p$ .

Compare how  $\rho_{FB}(\Delta y)$  changes for increasing  $\Delta y = 0, 1, 2, 3, \dots$ , with and without MI switched on (`PartonLevel:MI = on/off`) for minimum-bias events (`SoftQCD:minBias = on`).

- Higgs production can proceed through several different production processes. For the Standard Model Higgs some process switches are:

`HiggsSM:ffbar2H` for  $f\bar{f} \rightarrow H^0$  (f generic fermion, here mainly  $b\bar{b} \rightarrow H^0$ );

`HiggsSM:gg2H` for  $gg \rightarrow H^0$ ;

`HiggsSM:ffbar2HZ` for  $f\bar{f} \rightarrow H^0 Z^0$ ;

`HiggsSM:ffbar2HW` for  $f\bar{f} \rightarrow H^0 W^\pm$ ;

`HiggsSM:ff2Hff(t:ZZ)` for  $f\bar{f} \rightarrow H^0 f\bar{f}$  via  $Z^0 Z^0$  fusion;

`HiggsSM:ff2Hff(t:WW)` for  $f\bar{f} \rightarrow H^0 f\bar{f}$  via  $W^+ W^-$  fusion;



`HiggsSM:all` for all of the above (and some more).

Study the  $p_{\perp}$  and  $\eta$  spectrum of the Higgs in these processes, and compare.

- You can also vary the Higgs mass with a `25:m0 = ...` and switch off FSR/ISR/MI as above for top.
- $Z^0$  production to lowest order only involves one process, accessible with `WeakSingleBoson:ffbar2gmZ = on`. The problem here is that the process is  $f\bar{f} \rightarrow \gamma^*/Z^0$  with full  $\gamma^*/Z^0$  interference and so a significant enhancement at low masses. The combined particle is always classified with code 23, however. So generate events and study the  $\gamma^*/Z^0$  mass and  $p_{\perp}$  distributions. Then restrict to a more “ $Z^0$ -like” mass range with `PhaseSpace:mHatMin = 75.` and `PhaseSpace:mHatMax = 120.`
- Using your favourite jet cluster algorithm, study the number of jets found in association with the  $Z^0$  above. You can switch off  $Z^0$  decay with `23:mayDecay = no`. If you do not have a jet finder around, to begin with you can use the simple `CellJet` one that comes with PYTHIA, see the “Event Analysis” page in the online manual. Again check the importance of FSR/ISR/MI.

## A The Event Record

The event record is set up to store every step in the evolution from an initial low-multiplicity partonic process to a final high-multiplicity hadronic state, in the order that new particles are generated. The record is a vector of particles, that expands to fit the needs of the current event (plus some additional pieces of information not discussed here). Thus `event[i]` is the  $i$ ’th particle of the current event, and you may study its properties by using various `event[i].method()` possibilities.

The `event.list()` listing provides the main properties of each particles, by column:

- `no`, the index number of the particle ( $i$  above);
- `id`, the PDG particle identity code (method `id()`);
- `name`, a plaintext rendering of the particle name (method `name()`), within brackets for initial or intermediate particles and without for final-state ones;
- `status`, the reason why a new particle was added to the event record (method `status()`);
- `mothers` and `daughters`, documentation on the event history (methods `mother1()`, `mother2()`, `daughter1()` and `daughter2()`);
- `colours`, the colour flow of the process (methods `col()` and `acol()`);
- `p_x`, `p_y`, `p_z` and `e`, the components of the momentum four-vector  $(p_x, p_y, p_z, E)$ , in units of GeV with  $c = 1$  (methods `px()`, `py()`, `pz()` and `e()`);
- `m`, the mass, in units as above (method `m()`).

For a complete description of these and other particle properties (such as production and decay vertices, rapidity,  $p_{\perp}$ , etc), open the program’s online documentation in a browser (section 2 point 6) and scroll down to the “Study Output” section, the “Particle Properties” link in the left-hand-side manu. For brief summaries on the less trivial of the ones above, read on.

## A.1 Identity codes

A complete specification of the PDG codes is found in the Review of Particle Physics [3]. An online listing is available from

[http://pdg.lbl.gov/2008/mcdata/mc\\_particle\\_id\\_contents.html](http://pdg.lbl.gov/2008/mcdata/mc_particle_id_contents.html)

A short summary of the most common id codes would be

1	d	11	$e^-$	21	g	211	$\pi^+$	111	$\pi^0$	213	$\rho^+$	2112	n
2	u	12	$\nu_e$	22	$\gamma$	311	$K^0$	221	$\eta$	313	$K^{*0}$	2212	p
3	s	13	$\mu^-$	23	$Z^0$	321	$K^+$	331	$\eta'$	323	$K^{*+}$	3122	$\Lambda^0$
4	c	14	$\nu_\mu$	24	$W^+$	411	$D^+$	130	$K_L^0$	113	$\rho^0$	3112	$\Sigma^-$
5	b	15	$\tau^-$	25	$H^0$	421	$D^0$	310	$K_S^0$	223	$\omega$	3212	$\Sigma^0$
6	t	16	$\nu_\tau$			431	$D_s^+$			333	$\phi$	3222	$\Sigma^+$

Antiparticles to the above, where existing as separate entities, are given with a negative sign.

Note that simple meson and baryon codes are constructed from the constituent (anti)quark codes, with a final spin-state-counting digit  $2s + 1$  ( $K_L^0$  and  $K_S^0$  being exceptions), and with a set of further rules to make the codes unambiguous.

## A.2 Status codes

When a new particle is added to the event record, it is assigned a positive status code that describes why it has been added, as follows:

code range	explanation
11 – 19	beam particles
21 – 29	particles of the hardest subprocess
31 – 39	particles of subsequent subprocesses in multiple interactions
41 – 49	particles produced by initial-state-showers
51 – 59	particles produced by final-state-showers
61 – 69	particles produced by beam-remnant treatment
71 – 79	partons in preparation of hadronization process
81 – 89	primary hadrons produced by hadronization process
91 – 99	particles produced in decay process, or by Bose-Einstein effects

Whenever a particle is allowed to branch or decay further its status code is negated (but it is *never* removed from the event record), such that only particles in the final state remain with positive codes. (Thus the `isFinal()` method returns `true/false` for positive/negative status.) Note that the “same” particle can occur in several copies, e.g. when parton showering gives it a changed momentum, but only at most one of them with positive status.

## A.3 History information

The two mother and two daughter indices of each particle provide information on the history relationship between the different entries in the event record. The detailed rules depend on the particular physics step being described, as defined by the status code. As an example, in a  $2 \rightarrow 2$  process  $ab \rightarrow cd$ , the locations of  $a$  and  $b$  would set the mothers

of  $c$  and  $d$ , with the reverse relationship for daughters. When the two mother or daughter indices are not consecutive they define a range between the first and last entry, such as a string system consisting of several partons fragment into several hadrons.

There are also several special cases. One such is when “the same” particle appears as a second copy, e.g. because its momentum has been shifted by it taking a recoil in the dipole picture of parton showers. Then the original has both daughter indices pointing to the same particle, which in its turn has both mother pointers referring back to the original. Another special case is the description of ISR by backwards evolution, where the mother is constructed at a later stage than the daughter, and therefore appears below in the event listing.

If you get confused by the different special-case storage options, the two `event(i).motherList()` and `event(i).daughterList()` methods are able to return a `vector` of all mother or daughter indices.

## A.4 Colour flow information

The colour flow information is based on the Les Houches Accord convention [4]. In it the number of colours is assumed infinite, so that each new colour line can be assigned a new separate colour. These colours are given consecutive labels: 101, 102, 103, .... A gluon has both a colour and an anticolour label, an (anti)quark only (anti)colour.

While colours are traced consistently through hard processes and parton showers, the subsequent beam-remnant-handling step often involves a drastic change of colour labels. Firstly, previously unrelated colours and anticolours taken from the beams may at this stage be associated with each other, and be relabelled accordingly. Secondly, it appears that the close space-time overlap of many colour fields leads to reconnections, i.e. a swapping of colour labels, that tends to reduce the total length of field lines.

## References

- [1] T. Sjöstrand, S. Mrenna and P. Skands, *Comput. Phys. Comm.* **178** (2008) 852 [arXiv:0710.3820]
- [2] T. Sjöstrand, S. Mrenna and P. Skands, *JHEP* **05** (2006) 026 [hep-ph/0603175]
- [3] Particle Data Group, C. Amsler et al., *Physics Letters* **B667** (2008) 1
- [4] E. Boos et al., in the *Proceedings of the Workshop on Physics at TeV Colliders*, Les Houches, France, 21 May - 1 Jun 2001 [hep-ph/0109068]